

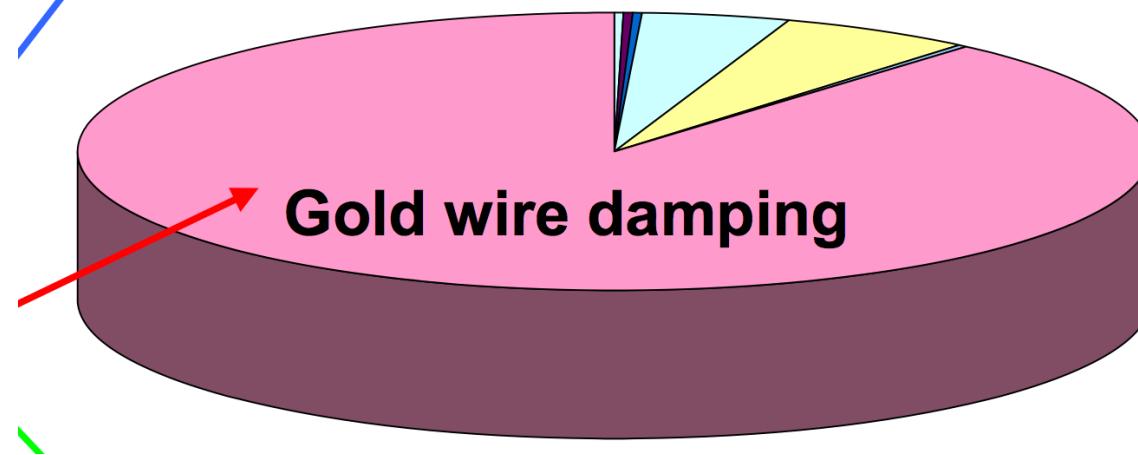
COLD DAMPING IN MICROSCOPE

A QUANTUM THERMODYNAMICAL ANALYSIS OF NOISE

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Noise/ Bias fluctuation



Total noise $= 4.89 \cdot 10^{-12} \text{ms}^{-2}/\text{Hz}^{1/2}$	Gold wire damping noise $= 4.85 \cdot 10^{-12} \text{ms}^{-2}/\text{Hz}^{1/2}$
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Main noise source : thermal noise of mechanical damping due to 7 μm gold wire

Active damping from servo control does not increase noise

The physics of cold damping

analysis in term of thermodynamics

ultimate limits and compatibility with quantum fluctuations

Using cold damping to improve measurements

early modelization of a capacitive accelerometer

Proposal for Microscope

Cold damping : Active reduction of thermal noise in electro- or opto-mechanical systems.

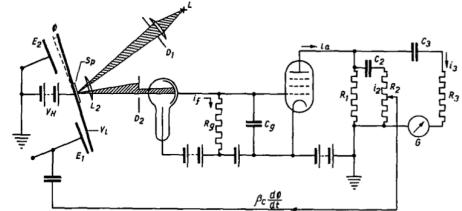
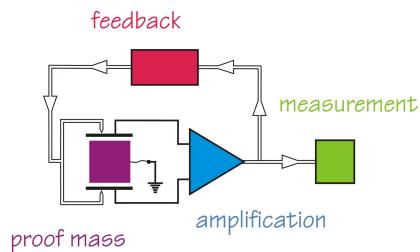
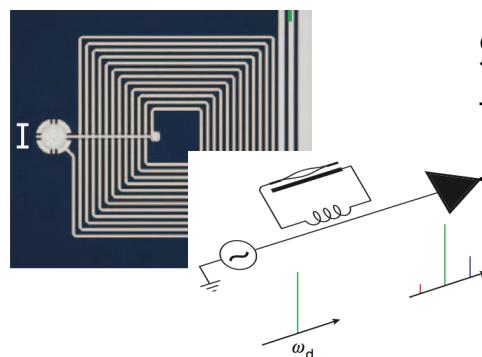


Fig. 1. Schematic arrangement of the electrometer, the damping circuit and the recording galvanometer.

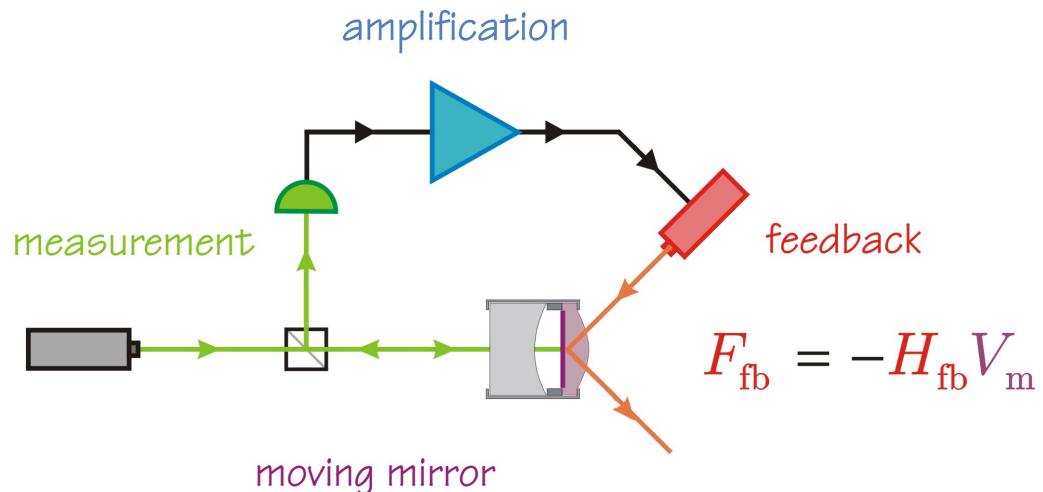
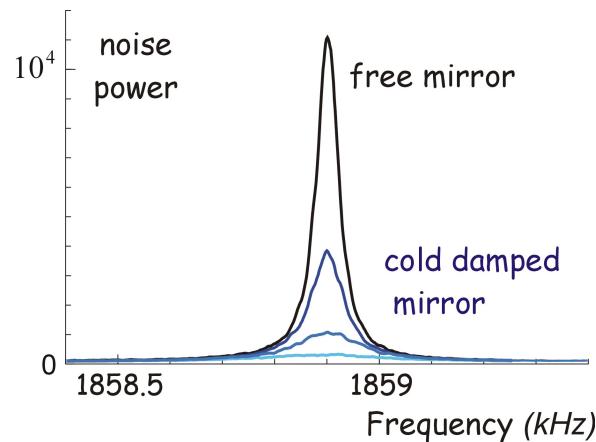
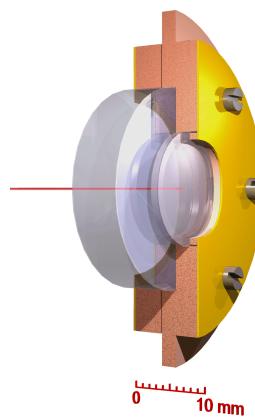
The reduction in the brownian motion of electrometers,
J.M.W. Milatz et al. Physica XIX, 195-207 (1953)



Quantum theory of fluctuations in a cold damped accelerometer.
F. Grassia, J.M. Courty, S. Reynaud, P. Touboul.
European Physical Journal D, 2000, 8, pp.101-110



Sideband cooling of micromechanical motion
to the quantum ground state
J. D. Teufel et al. Nature 475, 359–363 (21 July 2011)



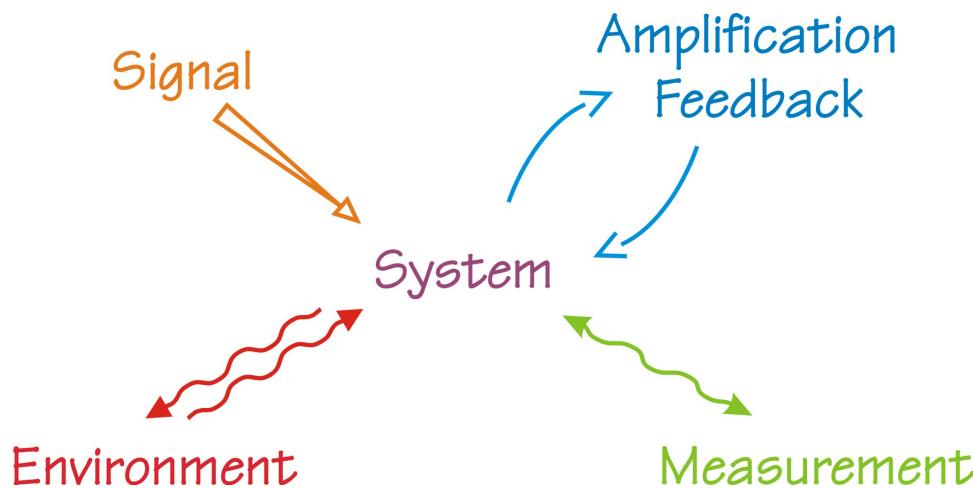
Mechanical oscillator is
an acoustic mode of a mirror

Motion is measured with Fabry Perot

Force is radiation pressure of a laser

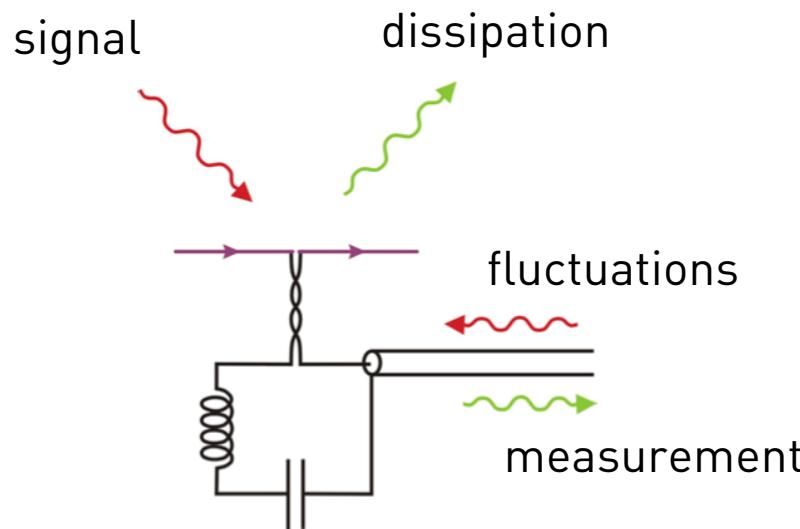
Cooling of a mirror by radiation pressure

P.F. Cohadon, A. Heidmann, M. Pinard, Phys. Rev. Lett. 83, 3174 (1999).

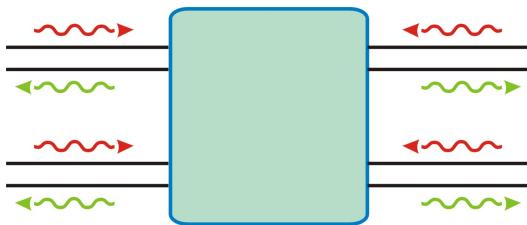


Important elements to consider in sensitivity analysis of actual measurements :

- Thermodynamical and quantum noises
- Active systems for signal amplification and servocontrols
- Spectral analysis of noise
- Modelization of complex devices

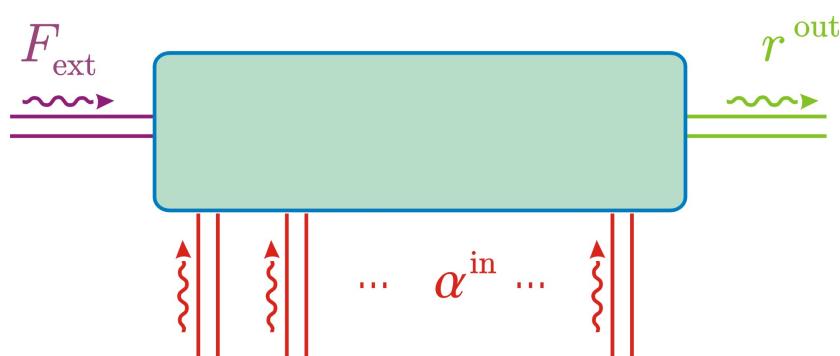


Fluctuation-dissipation theorem ensures consistency between the thermal noise of the oscillator and the coupled fluctuations bath



Unitarity of S matrix enforces
thermodynamic constraints
quantum constraints

Scattering of quantum fields



Force estimator

$$\hat{F}_{\text{ext}} \propto r^{\text{out}}$$

$$= F_{\text{ext}} + \sum_{\alpha} \mu_{\alpha} \alpha^{\text{in}}$$

Added noise

$$\Sigma_{FF} = \sum_{\alpha} |\mu_{\alpha}|^2 \sigma_{\alpha\alpha}^{\text{in}}$$

Thermal and quantum noises

$$\sigma_{aa}^{\text{in}}[\omega] = \frac{1}{2} \coth \frac{\hbar|\omega|}{k_B T_a}$$

$$= \frac{1}{\frac{\hbar|\omega|}{e^{k_B T_a} - 1}} + \frac{1}{2}$$



$$T \rightarrow \infty \quad \hbar |\omega| \sigma_{aa} \simeq k_B T_a$$

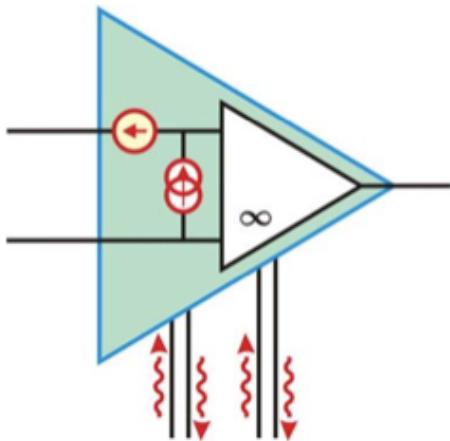
$$T \rightarrow 0 \quad \hbar |\omega| \sigma_{aa} \simeq \frac{1}{2} \hbar |\omega|$$

Amplification is also treated as a scattering.

Example : operational amplifier

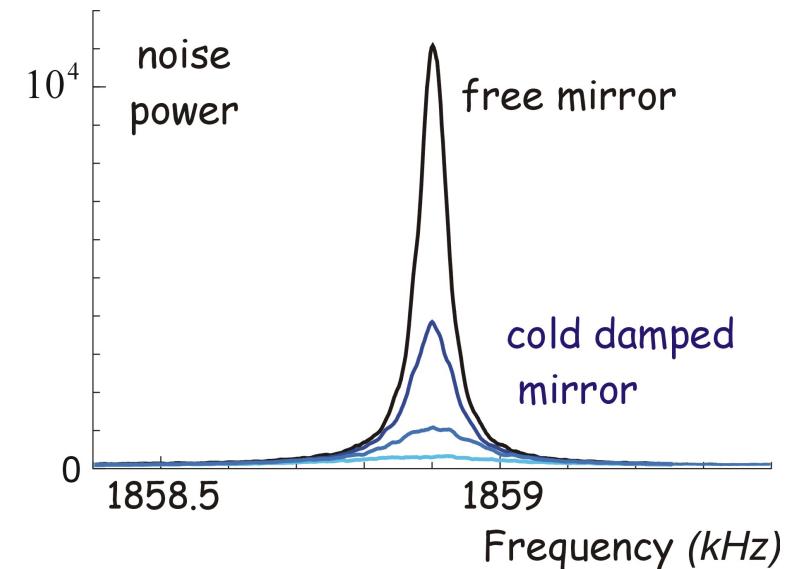
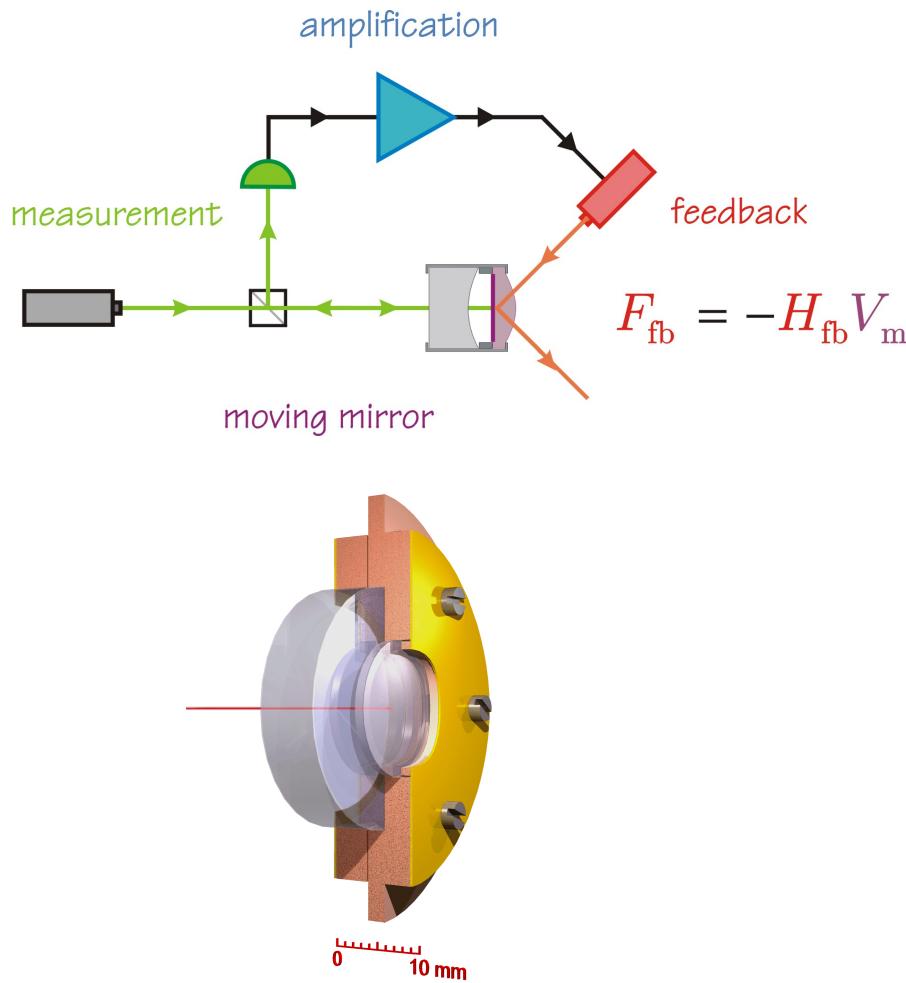
Voltage noise and current noise

Characterized by noise impedance and noise temperature



Charge and Flux are conjugated in quantum regime.

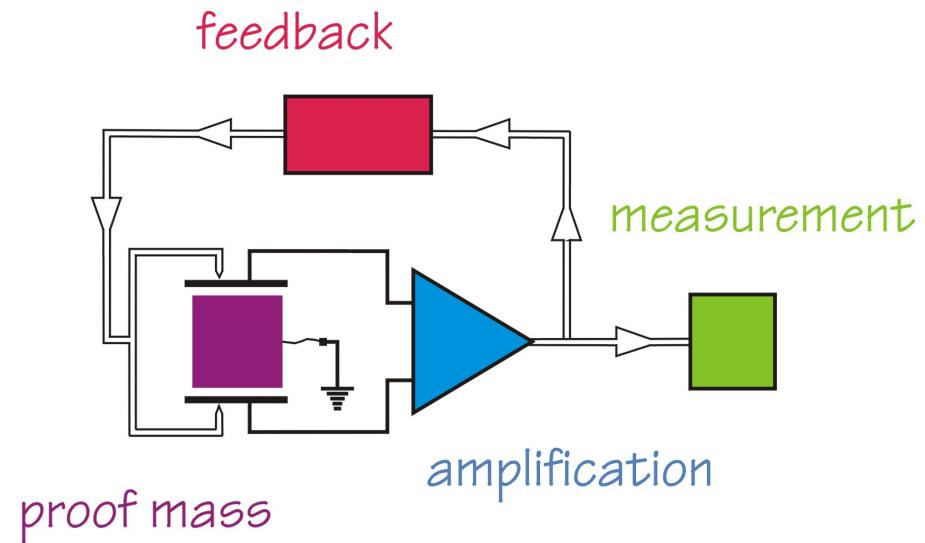
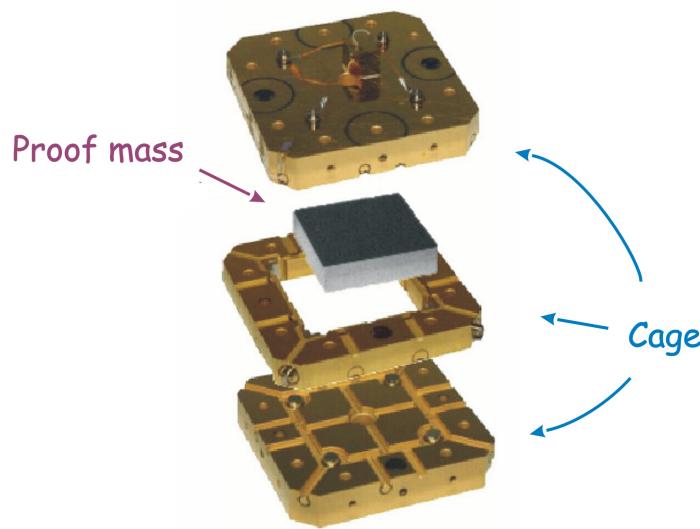
Quantum noise in ideal operational amplifiers,
Courtay, Grassia, Reynaud, Europhys. Lett. 46 (1999), 31-37



$$T = \frac{H_m}{H_m + H_{fb}} T_m$$

Cooling of a mirror by radiation pressure

P.F. Cohadon, A. Heidmann, M. Pinard, Phys. Rev. Lett. 83, 3174 (1999).



Detection
amplification noise
back action of servocontrol

Active control of proof mass
restoring force
cold damping

$$\sqrt{\Sigma_{aa}} = 1.2 \cdot 10^{-12} \text{ m s}^{-2} / \sqrt{\text{Hz}}$$

Quantum theory of fluctuations in a cold damped accelerometer.
F. Grassia, J.M. Courty, S. Reynaud, P. Touboul. E.P.J.I D, 2000, 8, pp.101-110

$$\hat{F}_{ext} = F_{ext} + \sum_{\alpha} \mu_{\alpha} \alpha^{\text{in}}$$

$$\mu_m = -\sqrt{2\hbar|\Omega|} \textcolor{red}{H}_m$$

$$\mu_{l_2} = -\frac{i\Omega\sqrt{\hbar}}{\sqrt{2}\textcolor{brown}{R}_l\omega_t} \Xi_m \quad \mu_{l_1} = 0$$

$$\mu_{r_1} = -\frac{\Omega\sqrt{\hbar}R_r}{2\sqrt{2\omega_t}Z_f} \Xi_m \quad \mu_{r_2} = 0$$

$$\mu_{a_1} = -\mu_{b_1} = \sqrt{2\hbar R_a \omega_t} \left(-\textcolor{brown}{x}_t + \frac{\Omega}{2\textcolor{brown}{x}_t \omega_t} \textcolor{violet}{Z}_f \Xi_m \right)$$

$$\mu_{a_2} = -\frac{i\Omega\sqrt{\hbar}R_a}{\sqrt{2}\textcolor{brown}{x}_t\sqrt{\omega_t}} \Xi_m \left(\frac{1}{R_a} - \frac{1}{R_l} - \frac{1}{Z_t} \right)$$

$$\mu_{b_2} = -\frac{i\Omega\sqrt{\hbar}R_a}{\sqrt{2}\textcolor{brown}{x}_t\sqrt{\omega_t}} \Xi_m \left(\frac{1}{R_a} + \frac{1}{R_l} + \frac{1}{Z_t} \right)$$

$$\Xi_m = \textcolor{red}{H}_m - iM\Omega + \frac{iK}{\Omega}$$

Characterization of all the measurement characteristics
measurement, noise, backaction, correlations

Noise of active control negligible
low temperature of amplifier noise
ratio of signal frequency and operation frequency

Sensitivity limited by residual damping

$$M = 0.27 \text{ kg}$$

$$H_m = 1.3 \cdot 10^{-5} \text{ kg s}^{-1}$$

$$\Theta_m = 300 \text{ K}$$

$$\Omega = 2\pi \cdot 5 \cdot 10^{-4} \text{ Hz}$$

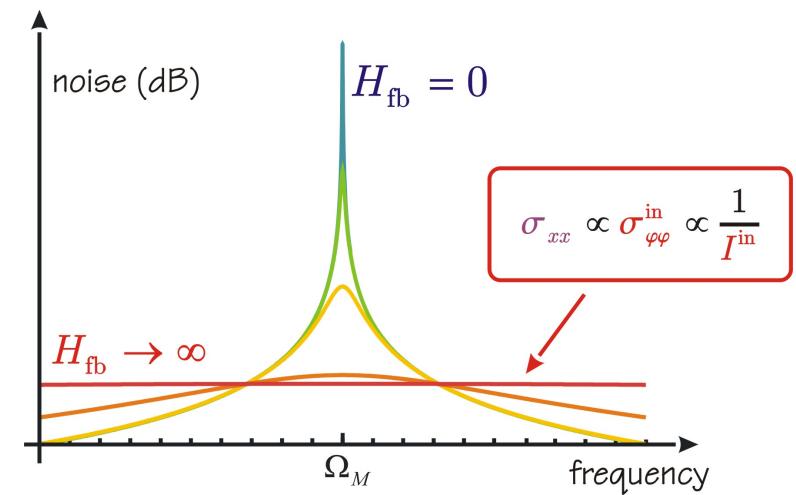
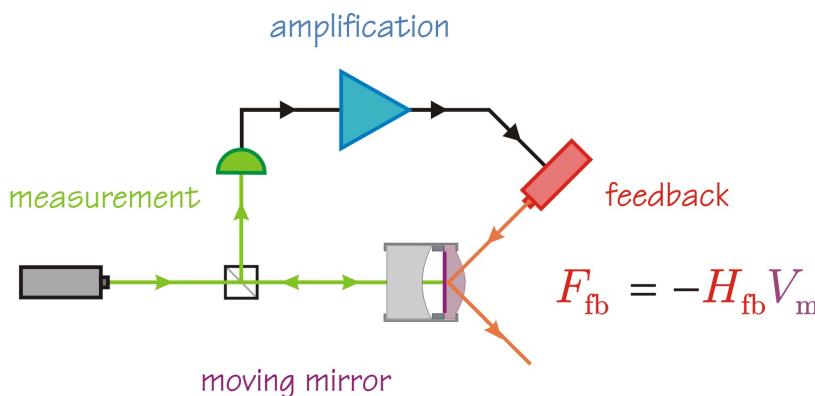
$$\omega_t = 2\pi \cdot 10^5 \text{ Hz}$$

$$R_a = 1.5 \cdot 10^5 \Omega$$

$$T_a = 1.5 \text{ K}$$

Quantum theory of fluctuations in a cold damped accelerometer.

F. Grassia, J.M. Courty, S. Reynaud, P. Touboul. E.P.J.D, 2000, 8, pp.101-110



$$E = \hbar \Omega_m \left(n_T \frac{H_m}{H_m + H_{fb}} + \frac{1}{2} \right)$$

Energy is reduced to ground state energy - \uparrow temperature is zero

J.M. Courty, A. Heidmann, M. Pinard Eur. Phys. J. D 17, 399 (2001)

Response to the call for ideas

Collaboration with Microscope Team in order to :

Perform the analysis on current sensor design

Evaluate effectiveness of cold damping

Use noise as a source of information on the instrument operation physics