

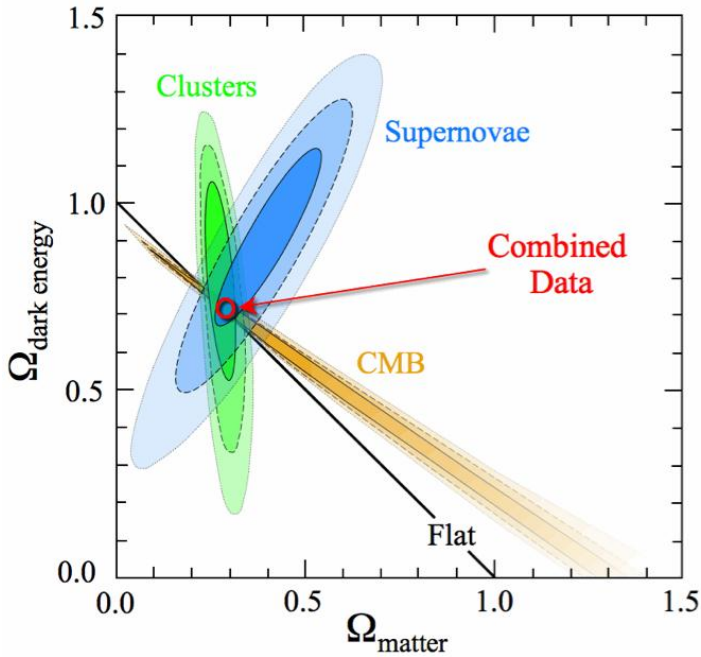
# Constraining screening mechanisms with MICROSCOPE

**Joel Bergé (ONERA)**  
**Jean-Philippe Uzan (IAP)**  
**Quentin Baghi (ONERA)**

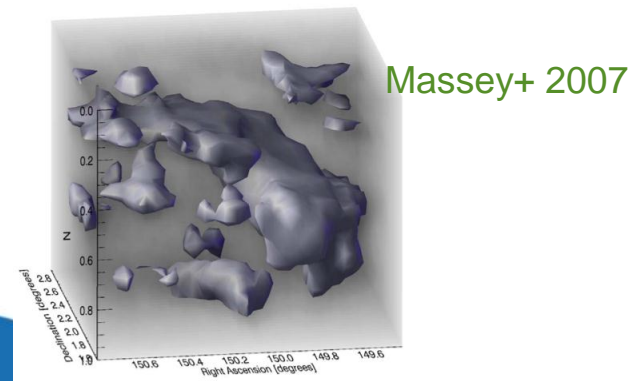


return on innovation

# Do we need new physics?

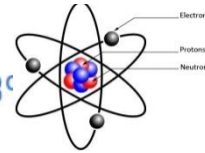
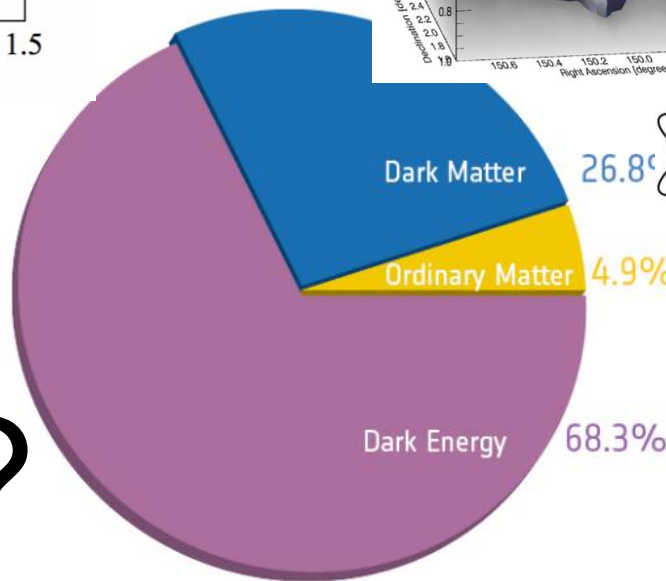


$\Lambda$ CDM model: explains acceleration of the Universe's expansion. But, is it the end of the road?



Dark sector...

?



...or new physics?

- Modified gravity
- Quantum gravity
- String theory
- New interaction

# New physics

- Modified gravity
  - Scalar-vector-tensor theories
  - Modified action theories (generalizes GR's action), eg  $f(R)$
  - ...
- Loop quantum gravity
- String theory
- Extra scalar field associated to a long range fifth force, coupled to matter (CDM and matter) -- chameleon, dilaton...
- ...

$$S_{\text{GR}} = \int \sqrt{-g} R d^4x$$

Predict Equivalence Principle violation e.g. due to coupling's dependence on matter species => finding such a violation will be a smoking gun for new physics beyond GR.

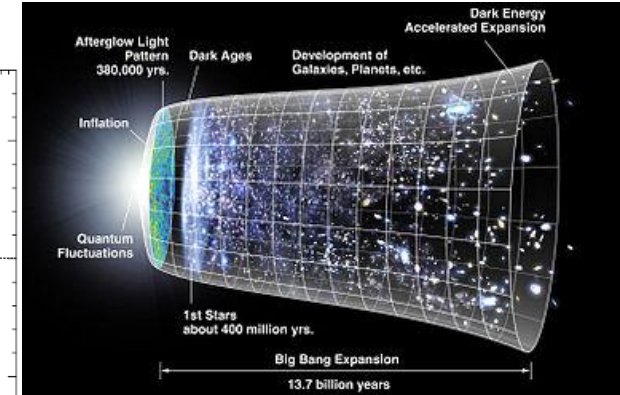
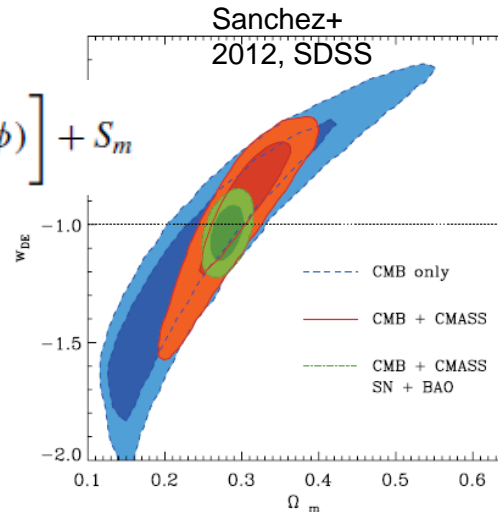


# Motivations for new scalar fields

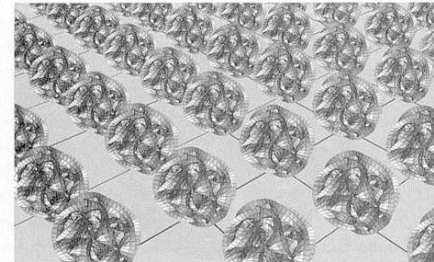
- Dark Energy, quintessence

$$S = \int d^4x \sqrt{-g} \left[ \frac{1}{2} M_{\text{pl}}^2 R - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right] + S_m$$

$$w \equiv \frac{P_\phi}{\rho_\phi} = \frac{\dot{\phi}^2/2 - V(\phi)}{\dot{\phi}^2/2 + V(\phi)}$$



- String theory: compactification into our low-energy, 4D space results in several massless scalar fields



According to string theory, the universe has extra dimensions curled up into a Calabi-Yau shape.

- Variation of constants

Some claims of variation of the fine structure in time of some parts in  $10^{15}$  (Uzan 2012)

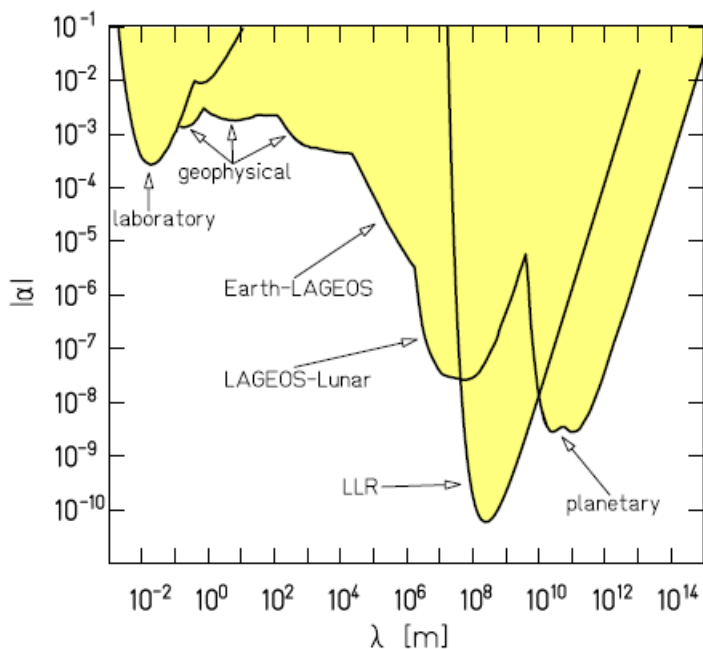
-> modeled as a coupling between matter and a scalar field

# The problem with massless scalar fields

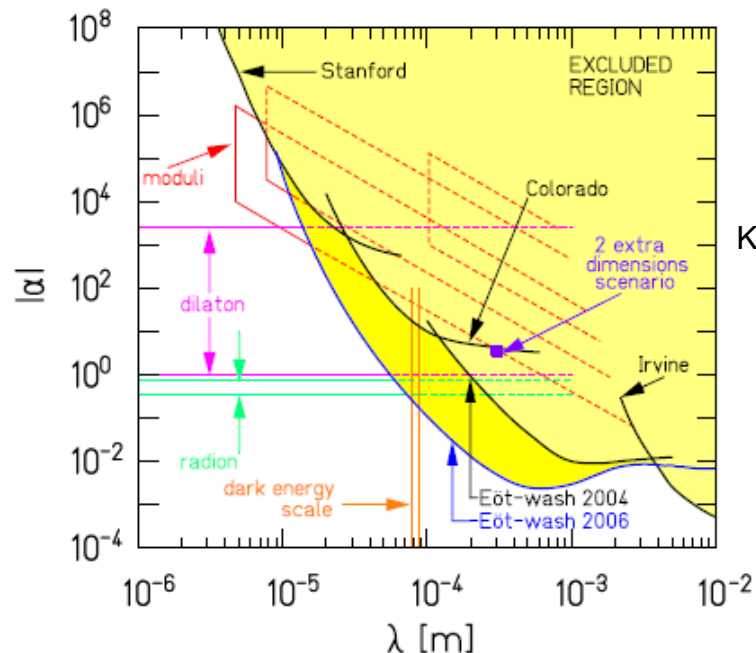
Long range => should be easily seen in Solar System / Earth experiments of  $1/r^2$  law and EP tests.

But we don't see them.

$$V(r) = -G \frac{m_1 m_2}{r} (1 + \alpha e^{-r/\lambda})$$



Gubser & Khoury 2004



Kapner+ 2006

*Don't they exist, or do they just hide themselves?*

Under some conditions, a scalar field which couples to matter can become hidden to our measurements and evade the constraints

⇒ The field has no detectable signature in these conditions, but behaves differently in other conditions. E.g., long-range in low-density regions (cosmological scales) but small-range in high-density regions (Earth, Solar System).

Zoology of screening mechanisms:

- Mass depends on local density: *chameleon*
- Coupling with matter depends on local density: *symmetron*, *Galileon*, *dilaton*
- Mass / coupling depends on local gravitational acceleration: *MOND*-type theories
- Coupling depends on local curvature: *Vainshtein* mechanism

# Chameleon in short (Khoury & Weltman 2004)

- Scalar field coupled to matter (with possibly different couplings between different matter species => can violate Equivalence Principle)
- Runaway potential, monotonic, decreasing
- Mass depends on local density
- Additional screening through thin-shell screening

## Abundant literature:

- Fifth force searches on Earth (Eöt-Wash)
- Solar System tests (Hees+ 2012)
- Cosmology (Brax+)





# Chameleon: more details (Khoury & Weltman 2004)

Action: 
$$S = \int d^4x \sqrt{-g} \left\{ \frac{M_{Pl}^2}{2} \mathcal{R} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right\} - \int d^4x \mathcal{L}_m(\psi_m^{(i)}, g_{\mu\nu}^{(i)})$$
  $\psi_m^{(i)}$ : matter fields

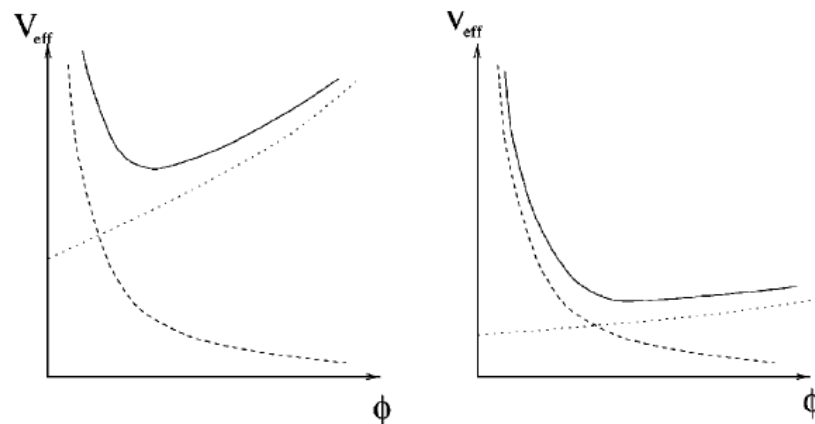
Potential  $V(\phi)$  of the runaway form. E.g Ratra-Peebles  $V(\phi) = M^{4+n} \phi^{-n}$ ,

Coupling to matter fields of the form  $e^{\beta_i \phi / M_{Pl}}$ ,  $\beta_i$ : dimensionless constants  $\sim 1$

Equation of motion 
$$\nabla^2 \phi = V_{,\phi} + \sum_i \frac{\beta_i}{M_{Pl}} \rho_i e^{\beta_i \phi / M_{Pl}}$$

=> dynamics of  $\phi$  are governed by the effective potential:

$$V_{eff}(\phi) \equiv V(\phi) + \sum_i \rho_i e^{\beta_i \phi / M_{Pl}}$$



Mass of the field: 
$$m_{min}^2 = V_{,\phi\phi}(\phi_{min}) + \sum_i \frac{\beta_i^2}{M_{Pl}^2} \rho_i e^{\beta_i \phi_{min} / M_{Pl}}$$

$$V_{,\phi}(\phi_{min}) + \sum_i \frac{\beta_i}{M_{Pl}} \rho_i e^{\beta_i \phi_{min} / M_{Pl}} = 0$$

$\phi_{min}$  and  $m_{min}$  depend on local density: larger  $\rho$  correspond to smaller  $\phi_{min}$  and larger mass => field can be massive enough on Earth to evade constraints but light enough in space to affect the gravitational dynamics (with no fine-tuning of  $\beta$ !).



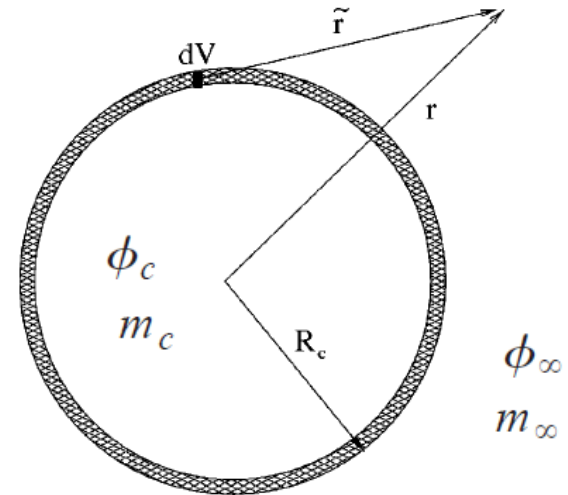
# Chameleon: profile and thin-shell screening

Goal: derive chameleon profile for a spherical compact object of mass  $M_c$ , radius  $R_c$  and density profile  $\rho(r)$ :

$$\rho(r) = \begin{cases} \rho_c & \text{for } r < R_c \\ \rho_\infty & \text{for } r > R_c \end{cases}$$

Equation of motion: 
$$\frac{d^2\phi}{dr^2} + \frac{2}{r} \frac{d\phi}{dr} = V_{,\phi} + \frac{\beta}{M_{Pl}} \rho(r) e^{\beta\phi/M_{Pl}}$$

with initial conditions 
$$\frac{d\phi}{dr} = 0 \quad \text{at } r=0, \quad \phi \rightarrow \phi_\infty \quad \text{as } r \rightarrow \infty$$



Inside the object,  $m_c \gg m_\infty$ ,  $\phi \sim \phi_c$ , a volume element  $dV$  contributes  $\exp(-m_\infty r) \Rightarrow$  exponentially suppressed. Only the volume elements close enough ( $\Delta R_c$ ) from the surface contribute to the exterior profile.

$$\phi(r) \approx - \left( \frac{\beta}{4\pi M_{Pl}} \right) \left( \frac{3\Delta R_c}{R_c} \right) \frac{M_c e^{-m_\infty r}}{r} + \phi_\infty \quad \frac{\Delta R_c}{R_c} = \frac{\phi_\infty - \phi_c}{6\beta M_{Pl} \Phi_c} \quad \text{assuming thin-shell condition } \left( \frac{\Delta R_c}{R_c} \right) \ll 1$$

For small objects,  $\frac{\Delta R_c}{R_c} > 1$  and 
$$\phi(r) \approx - \left( \frac{\beta}{4\pi M_{Pl}} \right) \frac{M_c e^{-m_\infty r}}{r} + \phi_\infty$$

No thin-shell screening

Thin-shell suppression factor

# Chameleon: fifth force, EP test and constraints

Chameleon force on a test particle of mass  $M$ :  $\vec{F}_\phi = -\frac{\beta}{M_{Pl}} M \vec{\nabla} \phi$

Profile on Earth + atmosphere (thin-shelled) and beyond:

$$\phi(r) \approx \begin{cases} \phi_\oplus & \text{for } 0 < r \leq R_\oplus, \\ \phi_{atm} & \text{for } R_\oplus \leq r \leq R_{atm}, \\ -\left(\frac{\beta}{4\pi M_{Pl}}\right) \left(\frac{3\Delta R_\oplus}{R_\oplus}\right) \frac{M_\oplus e^{-m_G(r-R_{atm})}}{r} + \phi_G & \text{for } r \geq R_{atm}, \end{cases} \quad \frac{\Delta R_\oplus}{R_\oplus} \equiv \frac{\phi_G - \phi_{atm}}{6\beta M_{Pl} \Phi_\oplus} < 10^{-7}$$

=> Fifth force on a test particle of mass  $M$  and coupling  $\beta_i$ :

$$|\vec{F}_\phi| = 2\beta\beta_i \left(\frac{3\Delta R_\oplus}{R_\oplus}\right) \frac{M_\oplus M}{8\pi M_{Pl}^2 r^2}$$

Magnitude of EP violation:  $\eta \equiv 2 \frac{|a_1 - a_2|}{a_1 + a_2} \sim 10^{-4} \beta^2 \frac{\Delta R_\oplus}{R_\oplus}$

Constraints on the chameleon-mediated interaction's range for a Ratra-Peebles potential

$$V(\phi) = M^{4+n} \phi^{-n}$$

Atmosphere  $m_{atm}^{-1} \leq 1 \text{ mm} - 1 \text{ cm}$ ,

Solar System  $m_G^{-1} \leq 10 - 10^4 \text{ AU}$ ,

Cosmological scales  $m_0^{-1} \leq 0.1 - 10^3 \text{ pc}$ ,

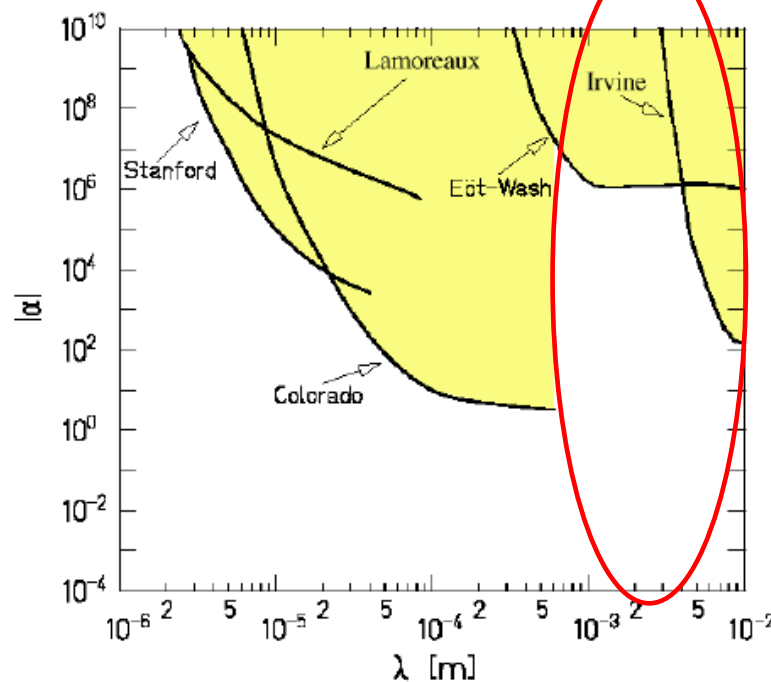
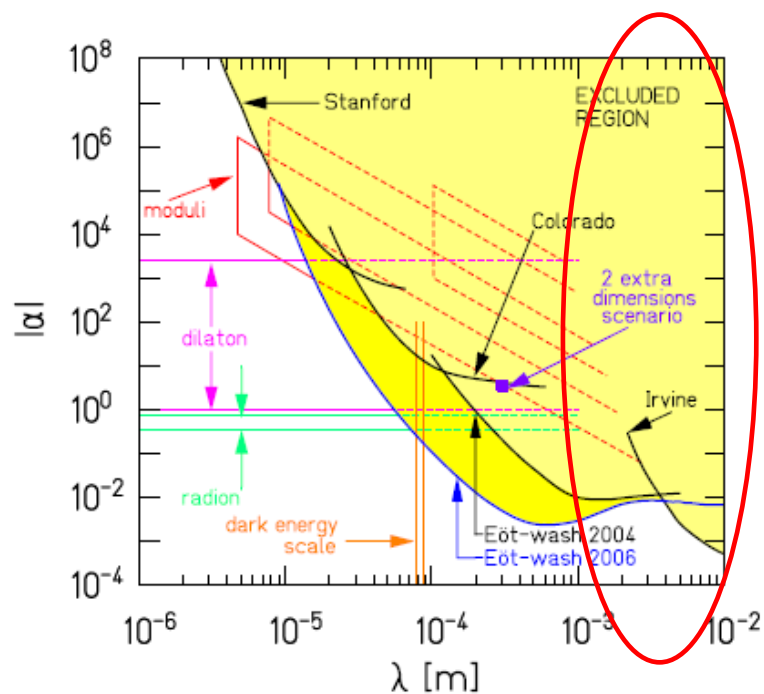
Behavior significantly different in space!

# Looser constraints on fifth force

Gubser & Khoury 2004

$$S = \int d^4x \sqrt{g} \left[ \frac{1}{2} (\partial\phi)^2 - \frac{1}{2} m_\phi^2 \phi^2 - \frac{\xi}{4!} \phi^4 \right] - \sum_\alpha \int \gamma_\alpha ds m_{i_\alpha}(\phi)$$

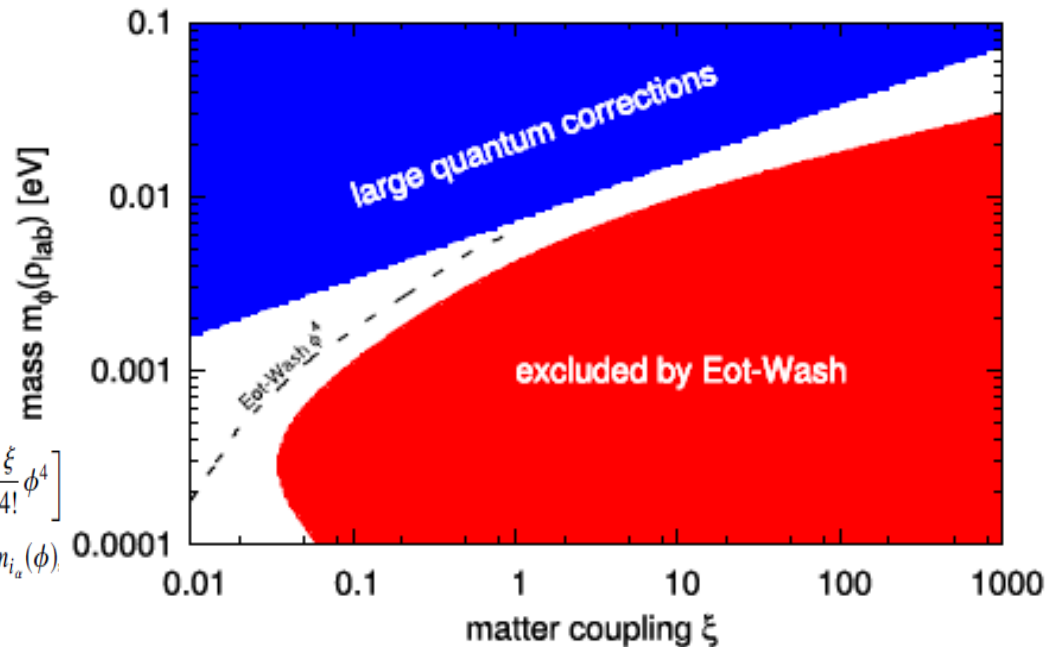
Thin-shelled



$$V(r) = -G \frac{m_1 m_2}{r} (1 + \alpha e^{-r/\lambda})$$

# Allowed mass and coupling values

Chameleon theories are effective field theories => quantum corrections should remain small compared to the classical potential => cannot have too large a mass



Upadhye+ 2012

Model-independent constraints from  $1/r^2$  law experiments

$$S = \int d^4x \sqrt{g} \left[ \frac{1}{2} (\partial\phi)^2 - \frac{1}{2} m_\phi^2 \phi^2 - \frac{\xi}{4!} \phi^4 \right] - \sum_\alpha \int \gamma_\alpha ds m_{i_\alpha}(\phi)$$

Chameleon fields already very much constrained: a small improvement in experiments could rule out all chameleon models



# Expectation for chameleon detection with MICROSCOPE

Order of magnitude estimate, based on Khoury & Weltman 2004

MICROSCOPE can see a chameleon-induced WEP violation if it is not thin-shelled, i.e. if  $\Delta R_{MIC}/R_{MIC} > 1$

Chameleon (the Earth is thin-shelled):

$$\phi(r) \approx \begin{cases} \phi_{\oplus} & \text{for } 0 < r \leq R_{\oplus}, \\ \phi_{atm} & \text{for } R_{\oplus} \leq r \leq R_{atm}, \\ -\left(\frac{\beta}{4\pi M_{Pl}}\right)\left(\frac{3\Delta R_{\oplus}}{R_{\oplus}}\right)\frac{M_{\oplus}e^{-m_G(r-R_{atm})}}{r} + \phi_G & \text{for } r \geq R_{atm}, \end{cases} \quad \frac{\Delta R_{\oplus}}{R_{\oplus}} \equiv \frac{\phi_G - \phi_{atm}}{6\beta M_{Pl}\Phi_{\oplus}} < 10^{-7}$$

At  $r=700\text{km}$ ,  $\phi(r) \sim \phi_G$

MICROSCOPE's Newtonian potential  $\sim 10^{-15}\Phi_{\oplus}$

$$\left. \begin{array}{l} \text{At } r=700\text{km}, \phi(r) \sim \phi_G \\ \text{MICROSCOPE's Newtonian} \\ \text{potential } \sim 10^{-15}\Phi_{\oplus} \end{array} \right\} \Delta R_{MIC}/R_{MIC} > 1 \text{ if } \frac{\Delta R_{\oplus}}{R_{\oplus}} > 10^{-15}$$

$\Rightarrow$  MICROSCOPE has no thin shell if

$$10^{-15} < \frac{\Delta R_{\oplus}}{R_{\oplus}} < 10^{-7}$$

$\Rightarrow$  EP violation  $\eta \approx 10^{-4}\beta^2 \frac{\Delta R_{\oplus}}{R_{\oplus}}$

$$\beta^2 \times 10^{-19} < \eta < \beta^2 \times 10^{-11}$$

## ***We need MICROSCOPE-specific predictions***

- Pick up our preferred screening mechanism(s)
- Derive trustworthy field equations in the satellite and precise expected physical effect on EP test.
- Link to full instrument (electronics and mechanics) simulator.
- Bricks already exist:
  - Simulink model of the instrument (performance group)
  - Physics simulation (OCA –G. Metris, L. Serron-- CMS)
  - Payload simulator at CNES

# The envisioned team

- Core members
  - Joel Bergé: ONERA Research scientist, member of MICROSCOPE CMS group, member of MICROSCOPE performance group
  - Jean-Philippe Uzan: IAP theoretical physicist
  - Quentin Baghi: ONERA PhD student
- A PhD student starting fall 2015?
  
- Performance group
- CMS
  
- Anyone interested

# Conclusion

- We have good reasons to add new scalar fields in physics
- To account for current tests of gravity, those scalar fields must either be very fine-tuned or remain hidden
- Several screening mechanisms have been proposed, that allow us to still add scalar fields
- EP violations are expected
- Significant EP violation (bigger than on Earth) could be seen with MICROSCOPE if a chameleon field exists.
- Otherwise, possibility to rule out all chameleons models.
- MICROSCOPE can be a unique experiment in the near future to make progress on constraining screening mechanisms.